Einstein's Summation Rule

Repeated indices are summed over

I have made a great discovery in mathematics!



inner product

$$a \cdot b = a_i b_i = \sum_i a_i b_i$$

Frobenius inner product

$$\langle A, B \rangle_F = A_{ij} B_{ij} = \sum_j \sum_i A_{ij} B_{ij}$$

Let's Practice Einstein's Summation Rule

$$tr(A) = ?$$

$$(A^T B)_{ij} = ?$$

$$tr(A^TB) = ?$$

$$\left(\vec{a}^T \vec{b} I\right)_{ij} = ?$$

$$a_i b_i = ?$$

$$a_{ii}a_{jj} = ?$$

$$a_{ii}a_{ii} = ?$$

Frobenius Inner Product $\langle A, B \rangle_F = A_{ij}B_{ij}$

$$A = \begin{bmatrix} \vec{a}_1, \vec{a}_2, \cdots, \vec{a}_n \end{bmatrix}$$

$$B = \begin{bmatrix} \vec{b}_1, \vec{b}_2, \cdots, \vec{b}_n \end{bmatrix}$$

$$\langle A, B \rangle_F = \sum_i \vec{a}_i^T \vec{b}_i$$

$$\langle RA, B \rangle_F = R_{ik} A_{kj} B_{ij} = \sum_i (R\vec{a}_i)^T \vec{b}_i$$

$$= A_{kj} (R^T B)_{kj} = \langle A, R^T B \rangle_F = \sum_i \vec{a}_i^T R \vec{b}_i$$

$$= R_{ik} (BA^T)_{ik} = \langle R, BA^T \rangle_F = \left\langle R, \sum_i \vec{b}_i \otimes \vec{a}_i \right\rangle_F$$

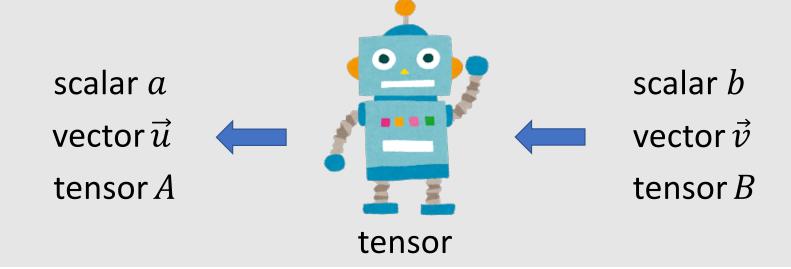
Tensor

What is Tensor?



In mathematics, a tensor is an algebraic object that describes a (multilinear) relationship between sets of algebraic objects related to a vector space.

https://en.wikipedia.org/wiki/Tensor

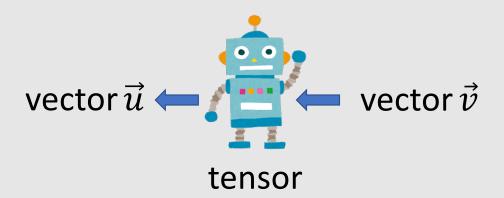


Two ways to Understand 2nd-order Tensor

Transformation by a tensor is give by the inner product

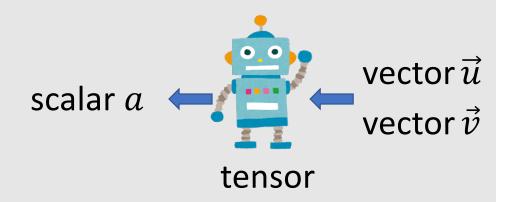
Linear form

$$\vec{u} = A \cdot \vec{v}$$



Quadratic form

$$a = \vec{u} \cdot (A \cdot \vec{v})$$



Outer Product (Tensor Product)

Outer product makes a tensor from two vectors

$$\vec{a} \otimes \vec{b} \qquad \qquad (\vec{a} \otimes \vec{b}) \cdot \vec{u} = \vec{a} (\vec{b} \cdot \vec{u})$$

• Tensor product $\vec{e} \otimes \vec{e}$ (||e|| = 1) defines projection

Definition

Projection ${\mathcal P}$

$$\mathcal{P}\big(\mathcal{P}(x)\big) = \mathcal{P}$$



Outer Product (Tensor Product)

Transformation for vectors in the outer product

check it out!

$$(A\vec{a})\otimes(B\vec{b})=?$$



Definition

$$\vec{a} \otimes \vec{b}$$

$$\vec{a} \otimes \vec{b}$$
$$(\vec{a} \otimes \vec{b}) \cdot \vec{u} = \vec{a}(\vec{b} \cdot \vec{u})$$

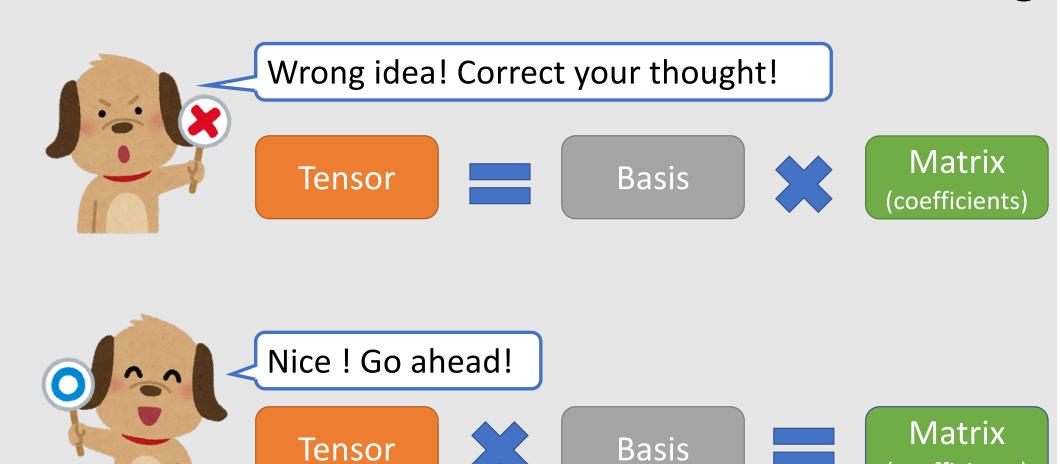
Tensor + Basis = Matrix

- Inner product with a basis vector gives a coefficient
 - This is true even if the basis is not orthonormal

$$v_i = \vec{v} \cdot \vec{e}_i$$

$$a_{ij} = \vec{e}_i \cdot (A \cdot \vec{e}_j)$$

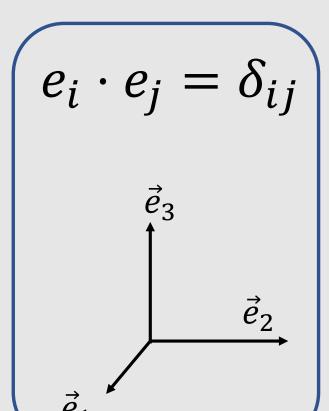
Tensor & Matrix: Common Misunderstanding



(coefficients)

Orthonormal Coordinates

Tensor can be written with bases and coefficients



$$v_{i} = \vec{v} \cdot \vec{e}_{i}$$

$$\Rightarrow \vec{v} = v_{i} \vec{e}_{i}$$

$$a_{ij} = \vec{e}_{i} \cdot (A \cdot \vec{e}_{j})$$

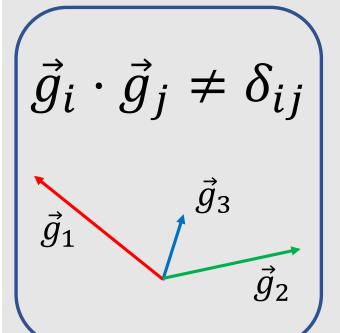
$$\Rightarrow A = a_{ij} \vec{e}_{i} \otimes \vec{e}_{j}$$



check it out!

Curvilinear Coordinates

- Non-orthogonal and un-normalized bases
- Dual bases solve the problem



dual basis

$$\vec{g}_i \cdot \vec{g}^j = \delta_i^j$$

$$\vec{g}^1$$
 \vec{g}^3 \vec{g}^2

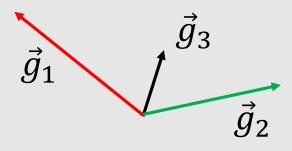
$$\vec{g}^1 = \frac{\vec{g}_2 \times \vec{g}_3}{\vec{g}_2 \cdot (\vec{g}_2 \times \vec{g}_3)}$$

$$\vec{g}^2 = \frac{\vec{g}_3 \times \vec{g}_1}{\vec{g}_3 \cdot (\vec{g}_3 \times \vec{g}_1)}$$

$$\vec{g}^3 = \frac{\vec{g}_1 \times \vec{g}_2}{\vec{g}_1 \cdot (\vec{g}_1 \times \vec{g}_2)}$$

Curvilinear Coordinates

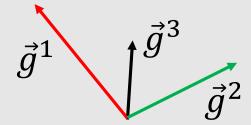
Expression of a vector in curvilinear coordinates



$$v^i = \vec{v} \cdot \vec{g}^i \qquad \qquad \vec{v} = v^i \vec{g}_i$$

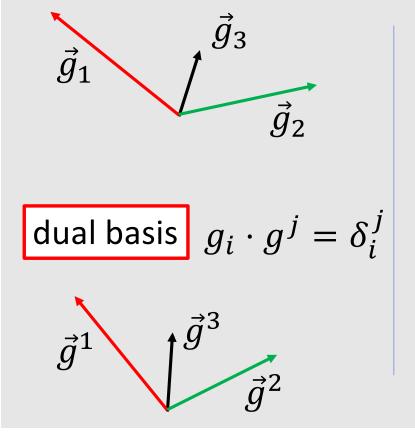
dual basis
$$g_i \cdot g^j = \delta_i^j$$

$$v_i = \vec{v} \cdot \vec{g}_i \qquad \qquad \vec{v} = v_i \vec{g}^i$$



Curvilinear Coordinates

Expression of a tensor in curvilinear coordinates



$$a^{ij} = \vec{g}^i \cdot (A \cdot \vec{g}^j) \longrightarrow A = a^{ij} \vec{g}_i \otimes \vec{g}_j$$

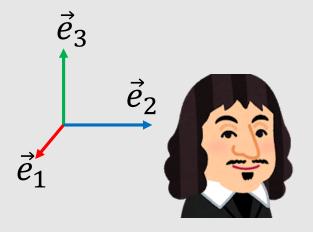
$$a_{ij} = \vec{g}_i \cdot (A \cdot \vec{g}_j) \longrightarrow A = a_{ij} \vec{g}^i \otimes \vec{g}^j$$

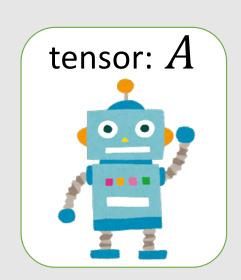
$$a^i_j = \vec{g}^i \cdot (A \cdot \vec{g}_i) \longrightarrow A = a^i_j \vec{g}_i \otimes \vec{g}^j$$

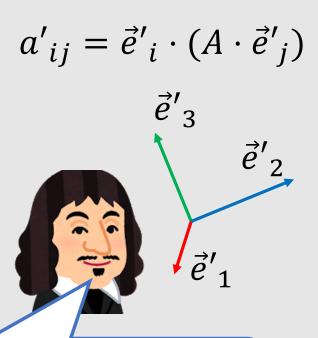
$$a^j_i = \vec{g}^i \cdot (A \cdot \vec{g}^j) \longrightarrow A = a^i_j \vec{g}^i \otimes \vec{g}^j$$

Coordinate Transformation

$$a_{ij} = \vec{e}_i \cdot (A \cdot \vec{e}_j)$$



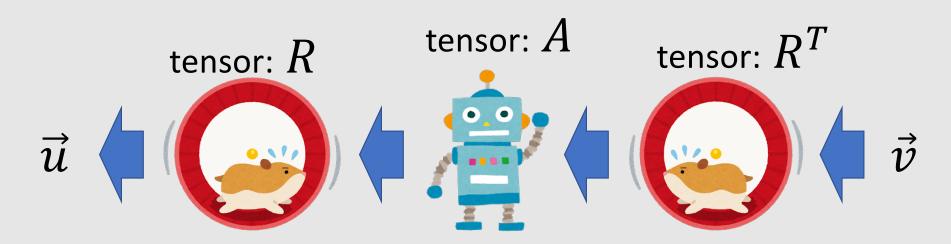




same tensor, different (coefficient) matrix!

Rotation of a Tensor

Un-rotating input and rotating output



$$\vec{u} = RAR^T \vec{v}$$

Rotation of a Tensor

Rotating bases while using the same coefficients

$$A = a_{ij}\vec{e}_i \otimes \vec{e}_j$$
rotation of basis
$$A' = a_{ij}(R\vec{e}_i) \otimes (R\vec{e}_j)$$

$$= a_{ij}R(\vec{e}_i \otimes \vec{e}_j)R^T$$

$$= R(a_{ij}\vec{e}_i \otimes \vec{e}_j)R^T$$

$$= RAR^T$$



check it out!

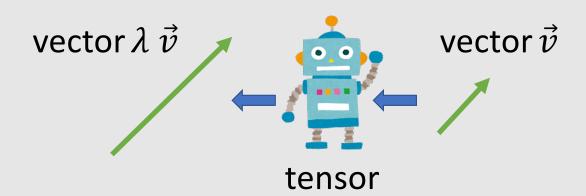
Simple Elastic Potential Energy for Continuum

Eigenvalue of Symmetric Tensor

• Eigenvalue of tensor is defined without matrix & coordinate

Linear form

$$\lambda \vec{v} = A \cdot \vec{v}$$



Eigenvalues and Frobenius Norm

from characteristic equation

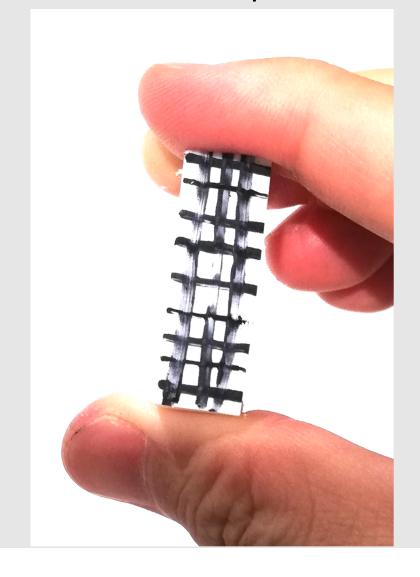
$$tr(A) = \lambda_1 + \lambda_2 + \lambda_3$$

$$A^2 \vec{e} = \lambda^2 \vec{e}$$

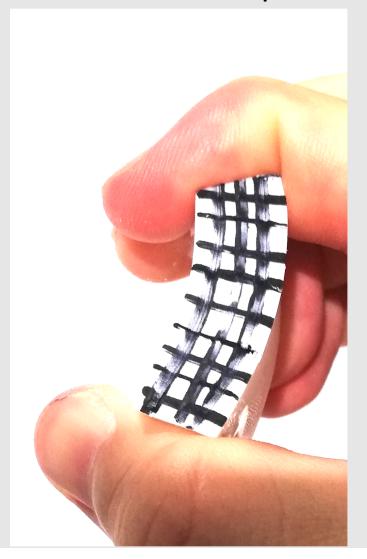
$$\operatorname{tr}(A^2) = \lambda_1^2 + \lambda_2^2 + \lambda_3^2$$

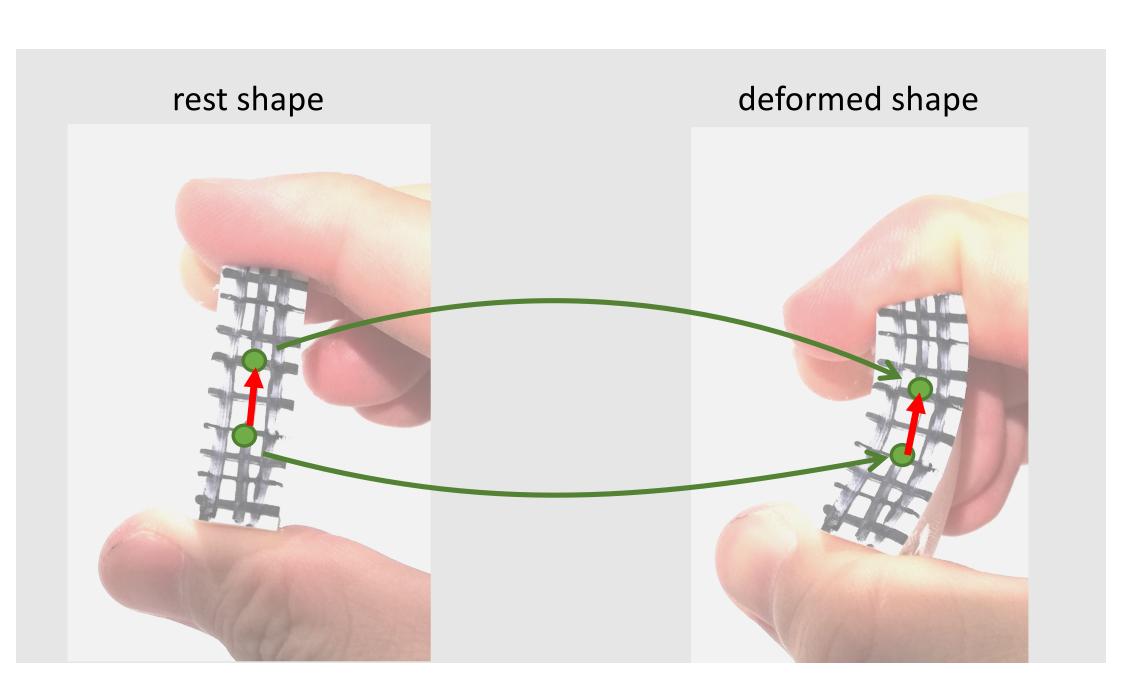
$$= \operatorname{tr}(A^T A) = \sum_{1 \le i,j \le 3} a_{ij}^2 = \left(\|A\|_F^2 \right)$$

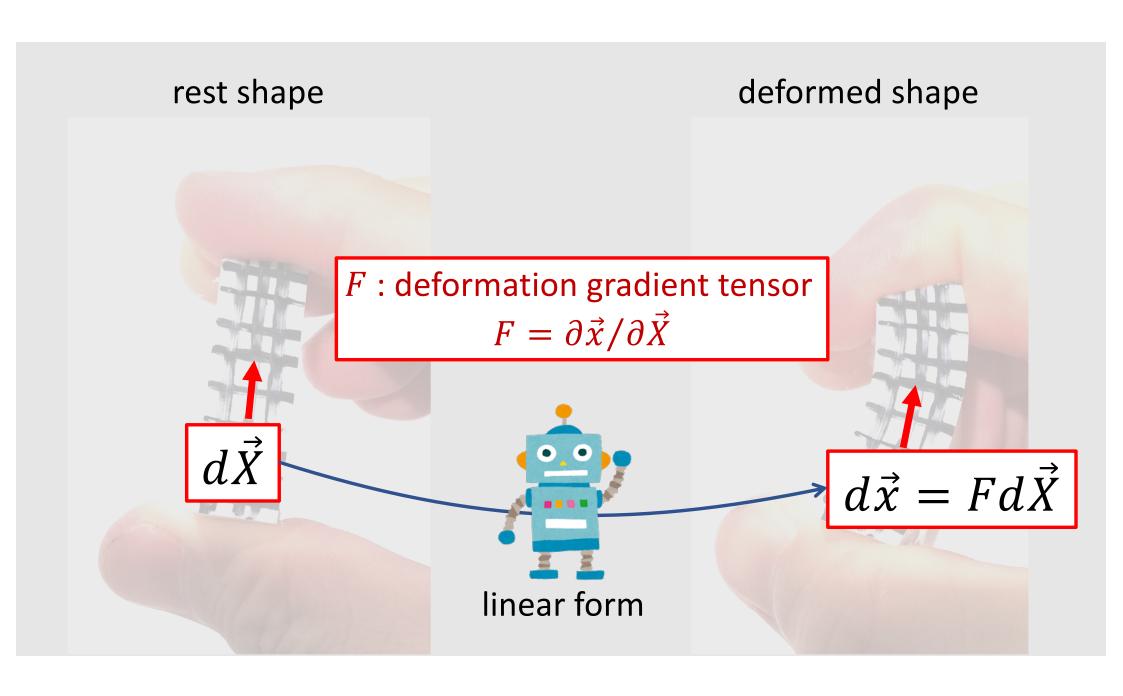
rest shape



deformed shape







SVD of Deformation Gradient Tensor

$$F = U\Sigma V^{T} \qquad \Sigma = \begin{bmatrix} \sigma_{1} & \\ & \sigma_{2} & \\ & & \\ \hline v^{T} & & \\ \hline v^{T} & & \\ \hline v_{2} & & \\ \hline v_{1} & & \\ \hline v^{T} & & \\ \hline v_{2} & & \\ \hline v_{1} & & \\ \hline v^{T} & & \\ \hline v_{2} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{2} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{2} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{2} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ \hline v_{1} & & \\ \hline v_{2} & & \\ \hline v_{3} & & \\ \hline v_{4} & & \\ \hline v_{5} & & \\ v_{5} & & \\ \hline v_{5} & & \\ v_{5} & & \\ \hline v_{5} & &$$

SVD of Deformation Gradient Tensor

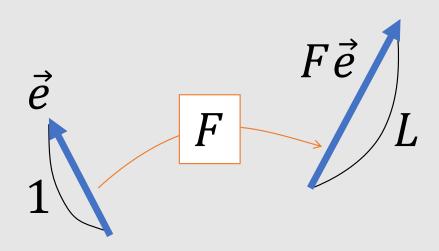
$$F = U\Sigma V^T$$
 $\Sigma = \begin{bmatrix} \sigma_1 & & \\ & \sigma_2 & \\ & & \sigma_3 \end{bmatrix}$

- Σ is the value we want for energy
- but SVD is costly
- How can we obtain Σ without SVD?

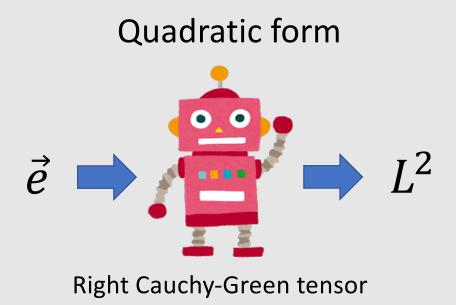


Gram Matrix F^TF Stands for Length Change

• $C = F^T F$: right Cauchy-Green tensor

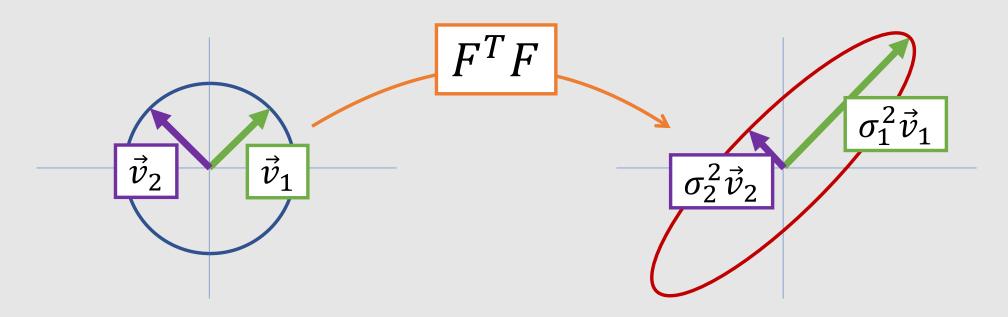


$$L^2(\vec{e}) = \vec{e}^T F^T F \vec{e}$$



Eigenvalue: Right Cauchy Green Tensor F^TF

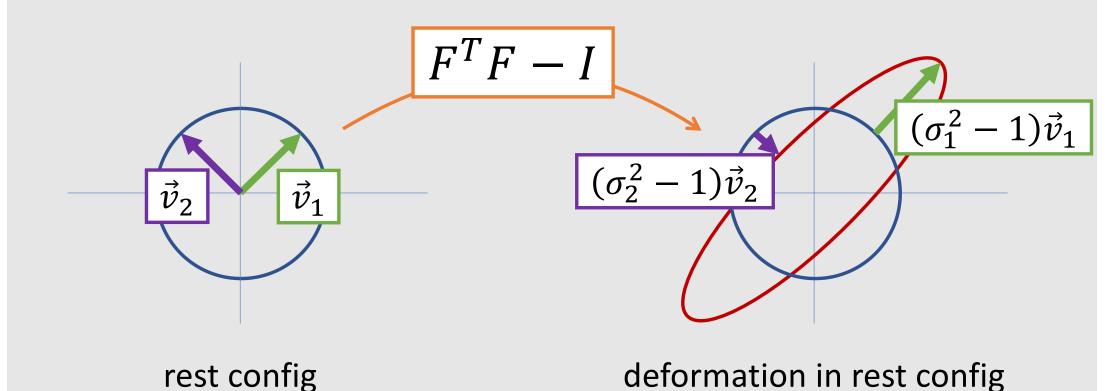
- Right Cauchy Green tensor is symmetric $F^TF = V\Sigma^2V^T$
- Eigenvalues of F^TF is squared of singular values : σ_1^2 , σ_2^2 , σ_3^2



rest config.

deformation in rest config.

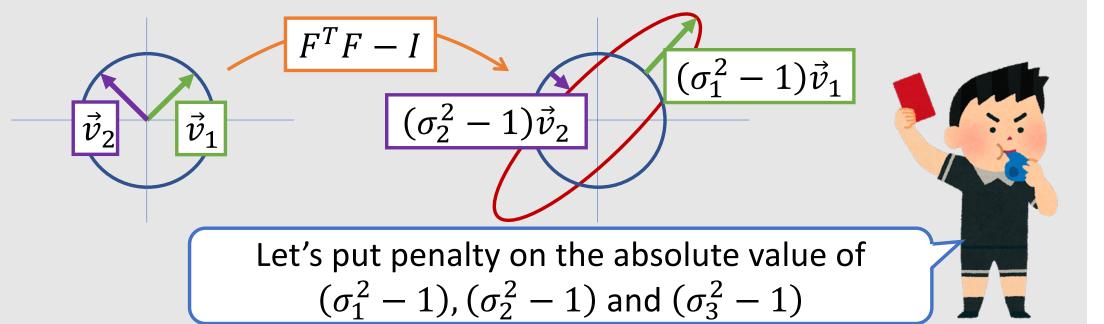
Eigenvalue: Green Lagrange Tensor $F^TF - I$



Making Energy from Eigenvalue

Energy for isotropic material

$$W(F) = ||F^T F - I||_F^2 = (\sigma_1^2 - 1)^2 + (\sigma_2^2 - 1)^2 + (\sigma_3^2 - 1)^2$$



Making Energy from Eigenvalue

Energy for isotropic material

$$W(F) = ||F^T F - I||_F^2 = (\sigma_1^2 - 1)^2 + (\sigma_2^2 - 1)^2 + (\sigma_3^2 - 1)^2$$



This energy doesn't have costly SVD and eigen decomposition easy to compute gradient & hessian!

Making Energy from Eigenvalue

$$W(F) = ||F^T F - I||_F^2 = (\sigma_1^2 - 1)^2 + (\sigma_2^2 - 1)^2 + (\sigma_3^2 - 1)^2$$

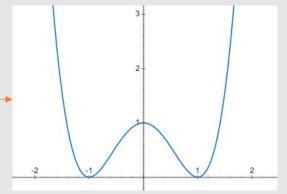
Wait... W(F) = 0 is not always no deformation. What about mirror reflection $\sigma_i = -1$?



$$y = (x^2 - 1)^2$$

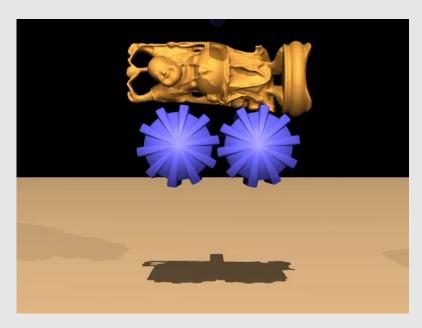


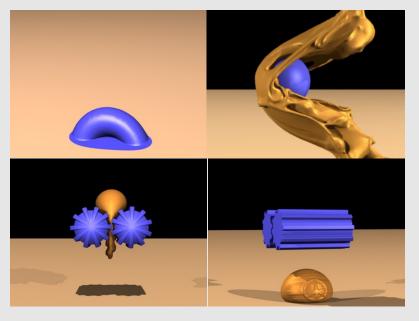
This energy is not robust to inversion



Invertible FEM [Irving et al. 2004]

• Elastic potential energy based on singular values of F that are σ_i , not on the eigen values of F^TF that are σ_i^2





G. Irving, J. Teran, and R. Fedkiw. 2004. Invertible finite elements for robust simulation of large deformation. In Proceedings of the 2004 ACM SIGGRAPH/Eurographics symposium on Computer animation (SCA '04)

Finite Element Method

What is Finite Element Method?

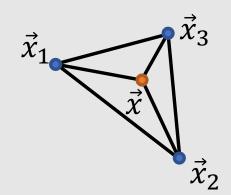
Solution by energy minimization

$$\vec{x}_{solution} = \underset{\vec{x}}{\operatorname{argmin}} W(\vec{x})$$



Value inside element is interpolated

$$\vec{x} = \sum_{i \in Nodes} w_i \vec{x}_i$$



Energy is integrated inside element and summed

$$W(\vec{x}) = \sum_{e \in Elements} W_e(\vec{x})$$

FEM of Laplace Equation on Triangle

Discrete Laplacian

the energy is sum of the squared differences between neighbors

Continuous Laplacian \longrightarrow the energy is integration of the squared gradient

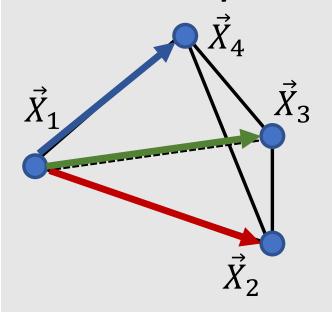
$$W(\phi) = \int_{\Omega} \nabla \phi \cdot \nabla \phi \ d\Omega$$

$$W_e(\phi) = \int_{\vec{x} \in Tri} \frac{\partial L_a \phi_a}{\partial \vec{x}} \cdot \frac{\partial L_b \phi_b}{\partial \vec{x}} \, d\vec{x} = \phi_a \phi_b \int_{\vec{x} \in Tri} \frac{\partial L_a}{\partial \vec{x}} \cdot \frac{\partial L_b}{\partial \vec{x}} \, d\vec{x}$$

Making the solution as smooth as possible!

Deformation Gradient Tensor F for Tet.

rest shape

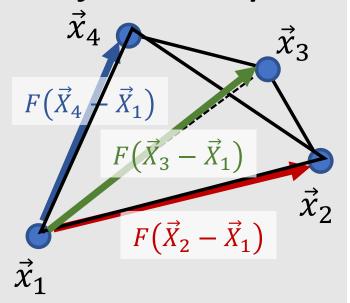


$$\vec{X}_3$$
 $(\vec{x}_2 - \vec{x}_1) = F(\vec{X}_2 - \vec{X}_1)$

$$(\vec{x}_3 - \vec{x}_1) = F(\vec{X}_3 - \vec{X}_1)$$

$$(\vec{x}_4 - \vec{x}_1) = F(\vec{X}_4 - \vec{X}_1)$$

deformed shape



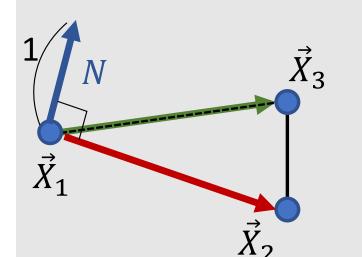
$$[\vec{x}_2 - \vec{x}_1, \vec{x}_3 - \vec{x}_1, \vec{x}_4 - \vec{x}_1] = F[\vec{X}_2 - \vec{X}_1, \vec{X}_3 - \vec{X}_1, \vec{X}_4 - \vec{X}_1]$$

$$F = [\vec{x}_2 - \vec{x}_1, \vec{x}_3 - \vec{x}_1, \vec{x}_4 - \vec{x}_1] [\vec{X}_2 - \vec{X}_1, \vec{X}_3 - \vec{X}_1, \vec{X}_4 - \vec{X}_1]^{-1}$$

Deformation Gradient Tensor F for 3D Tri.

rest shape

$(\vec{x}_2 - \vec{x}_1) = F(\vec{X}_2 - \vec{X}_1)$



$$\vec{X}_3$$
 $(\vec{x}_3 - \vec{x}_1) = F(\vec{X}_3 - \vec{X}_1)$

$$n = FN$$

$$F(\vec{X}_3 - \vec{X}_1)$$

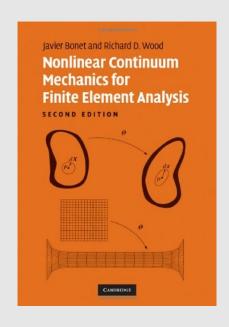
$$\vec{X}_1$$

$$\vec{X}_1$$

$$[\vec{x}_2 - \vec{x}_1, \vec{x}_3 - \vec{x}_1, n] = F[\vec{X}_2 - \vec{X}_1, \vec{X}_3 - \vec{X}_1, N]$$

$$F = [\vec{x}_2 - \vec{x}_1, \vec{x}_3 - \vec{x}_1, n] [\vec{X}_2 - \vec{X}_1, \vec{X}_3 - \vec{X}_1, N]^{-1}$$

Reference



•Bonet, Javier, and Richard D. Wood. 1997. *Nonlinear continuum mechanics for finite element analysis*. Cambridge: Cambridge University Press.

Linearization of Green-Lagrange Strain

Green Lagrange Strain: $E = F^T F - I$

More Realistic Elasticity Model

Green Lagrange Strain: $E = F^T F - I$

$$W(F) = ||F^T F - I||_F^2$$

$$W(F) = ||F^T F - I||_F^2$$

4th-order Tensor

Definition of Inner product in higher dimension

